

Figure 1. Perspective drawing of  $[Li(bhppH_2)]^-$  showing the molecular structure and labeling scheme.

Positional parameters for the non-hydrogen atoms are given in Table I; selected bond distances and angles, in Table II; structure factors, positional parameters for hydrogen atoms, and anisotropic thermal parameters are available as supplementary material. The molecular geometry and labeling diagram are shown in Figure 1.

### **Results and Discussion**

The interaction of  $Zn(bhppH_2)_2 \cdot 2CH_3OH$  with 6 equiv of LiMeO in ethanol at 0 °C produces a yellowish homogeneous solution that, upon the addition of NEt<sub>4</sub>Cl and a large excess of water, yields an amorphous solid. Purification of the solid was achieved by treatment with warm ethyl acetate. The insoluble residue was discarded. Cooling the warm solution to room temperature afforded yellow crystals that were suitable for X-ray diffraction studies. The compound is air stable.

Until recently, few well-characterized mononuclear lithium complexes where the only metal present is Li were known. Some of these complexes have at least one C-Li bond, examples are 2-lithio-2-phenyl-1,3-dithiane-tetrahydrofuran-tetramethylethylenediamine<sup>10</sup> and Li[CH(SiMe<sub>3</sub>)<sub>2</sub>](pmdeta).<sup>11</sup> Others are not organometallic, but in them, lithium reaches tetrahedral coordination by means of different kind of ligands in a tridimensional arrangement. In most of them hydrogen bonds are involved. One example is lithium hydrogen acetylenedicarboxylate hydrate.<sup>12</sup> In this case the lithium is surrounded by four oxygen atoms from three anions and one water molecule. Another example is  $(B15C5)Li(picrate)\cdot 2H_2O^{13}$  (B15C5 benzo-15-crown-5). In it the cation is coordinated by two oxy atoms from the chelating picrate anion and two water molecules Surprisingly, the crown ether molecule is not bound to the Linea but to water molecules through weak hydrogen bonds. To one knowledge there is no structural information about a discrete molecular complex as simple as the one presented in this paper, where the tetrahedral coordination around the Li ion is achieved by donor atoms from two identical bidentate  $(bhppH_2)$  molecules. There was more information about these kind of compounds in solution where a few complexes of the type 1:2 Li:L  $6ag_{in}$  L = dipivaloylethane, diisobutyrylmethane)<sup>14</sup> had been described.

The structure of NEt<sub>4</sub>[Li(bhppH<sub>1</sub>) as revealed by X-ray crystallography shows a Li atom surrounded by four oxygen atoms in a distorted tetrahedral geometry where the 1,3-bis(2hydroxyphenyl)-1,3-propanedionate ligand acts in a bidentate

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fashion. The Li-propanedionate rings adopt a half-chair form with Li deviating 0.10 (2) and 0.31 (2) Å from the mean plane defined by the propanedionate moiety. The most planar ring shows the shorter Li-O bond distances while the dihedral angle between the two ring planes is equal to 99 (1)°. The phenyl rings linked to C(1), C(3), C(4), and C(6) atoms are twisted by 10, 11, 12, and 13°, respectively with regard to the plane defined by the  $\beta$ -diketone moiety. These twistings and the C(sp<sup>2</sup>)-C(phenyl) bond distances (e.g., C(1)-C(11)) indicate that the electron delocalization of the propanedionate ring does not extend to the phenyl rings. The hydroxy groups are hydrogen bonded to the adjacent metal-coordinated O atoms (O(1)-O(12), O(3)-O(32), O(4)...O(42), and O(6)...O(62) with distances equal to 2.47 (1), 2.47 (1), 2.44 (1), and 2.48 (1) Å, respectively). The Li-O bond distances are all different with values 1.85(2), 1.90(2), 1.89(2), and 1.91 (2) Å (average value 1.885 (5) Å). This large Li-O bond distance is very similar to the value 1.99 (9) Å (average) found in LiSiON $\alpha^{12}$  or 1.97 (5) Å found in 2-lithio-2-phenyl-1,3-dithiane-tetrahydrofuran-tetramethylethylene, precluding, consequently, the possibility of aromatic character in the six-membered  $\beta$ -diketonate chelate ring. On the other hand, the C++O distances of 1.29 (4) Å (average) found in this NEt<sub>4</sub>[Li(bhppH<sub>2</sub>)<sub>2</sub>] complex are similar to these observed in  $Zn(bhppH_2) \cdot 2EtOH$  (C=O = 1.303 (13) Å). This close similarity is a clear indication of the lack of aromatic character in the  $Zn-\beta$ -diketone ring of Zn-(bhppH<sub>2</sub>)<sub>2</sub>·2EtOH.

The interaction of  $Zn(bhppH_2)_2$ ·2EtOH with 6 equiv of Na-MeO, in conditions similar to these described to get  $NEt_4[Li-(bhppH_2)_2]$  yielded Na(bhppH\_2)·H<sub>2</sub>O, as evidenced by analysis and <sup>1</sup>H NMR.

Following a similar procedure, reaction of  $Zn(acac)_2H_2O$  with 6 equiv of LiMeO did not yield  $NEt_4[Li(acac)_2]$ . No further attempts were made to identify the nature of the resulting compounds.

These results in the solid state and others previously reported in solution for dipivaloylethane and diisobutyrylmethane<sup>14</sup> as complexing  $\beta$ -diketones show the great affinity of chelating agents of the type R"COCH(R')COR", where R', R" are bulky terminal groups, for the lithium ion. The affinity for Na<sup>+</sup> had been proven to decrease in solution for large  $\beta$ -diketones; e.g., the chelating tendency of dipivaloylethane for lithium was 1000 times greater than that for sodium.<sup>15</sup> As indicated earlier these results are in agreement with ours where [Li(bhppH<sub>2</sub>)<sub>2</sub>]<sup>-</sup> has been obtained but not [Na(bhppH<sub>2</sub>)<sub>2</sub>]<sup>-</sup>.

**Registry No.** NEt<sub>4</sub>[Li(bhppH<sub>2</sub>)<sub>2</sub>], 96212-06-5; Zn(bhppH<sub>2</sub>)<sub>2</sub>, 93401-40-2; Na(bhppH<sub>2</sub>), 96212-07-6; Zn(acac)<sub>2</sub>, 14024-63-6; LiMeO, 865-34-9.

Supplementary Material Available: Tables of structure factors, positional parameters for hydrogen atoms, and anisotropic thermal parameters (12 pages). Ordering information is given on any current masthead page.

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Symmetry Lowering in B<sub>3</sub>H<sub>9</sub> and B<sub>4</sub>H<sub>12</sub>

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The usual symmetry assignments are  $D_{3h}$  for  $B_3H_9$  and  $D_{4h}$  for  $B_4H_{12}$ , both of which are ring structures in which BH<sub>2</sub> groups are

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Table I. Absolute Energies (-hartrees) and Zero-Point Energies (kcal/mol) for B<sub>3</sub>H<sub>9</sub> and B<sub>4</sub>H<sub>12</sub> Optimized in Different Point Groups

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	3-21G	6-31G	MP2/6-31G	MP3/6-31G	6-31G*	ZPE <sup>a</sup>	
$B_{3}H_{9}(1), D_{3h}$	78.71636	79.13092	79.33843	79.37471	79.178 77	60.66 (3)	
$B_{3}H_{9}(2), C_{3}$	78.71780	79.132 40	79.342 89	79.37971	79.18208	62.44 (0)	
$B_4H_{12}$ (3), $D_{4h}$	104.95111	105.497 48	105.76428	105.81217	105.548 98	79.18 (5)	
$B_4H_{12}$ (4), $D_{2d}$	104.952 30	105.498 79	105.768 55	105.816 53	105.551 97	79.69 (4)	
$B_4H_{12}$ (5), $C_{4n}$	104.953 36	105.499 20	105.770 33	105.818 48	105.55627	82.31 (1)	
$B_4H_{12}$ (5), $C_{2v}$	104.953 97	105.500 53	105.77373	105.821 92	105.55813	82.21 (0)	

<sup>a</sup> Zero-point energy; the number in parentheses refers to the number of calculated negative (imaginary) frequencies.

Table II. Energies (kcal/mol) Relative to either  $B_3H_9$  or  $B_4H_{12}$  at Difference Levels of Theory<sup>a</sup>

	3-21G	6-31G	MP2/ 6-31G	MP3/ 6-31G	6-31G*	[MP3/ 6-31G*]	[MP3/6-31G*] + ZPC
$B_{3}H_{9}(1), D_{3k}$	0	0	0	0	0	0	0
$B_{3}H_{9}(2), C_{3v}$	-0.9	-0.9	-2.8	-2.8	-2.1	-4.0	-2.2
$B_4H_{12}$ (3), $D_{4h}$	0	0	0	0	0	0	0
$B_4H_{12}$ (4), $D_{2d}$	-0.7	-0.8	-2.7	-2.7	-1.9	-3.8	-3.3
$B_4H_{12}$ (5), $C_{4n}$	-1.4	-1.1	-4.0	-4.6	-4.2	-7.7	-4.6
$B_4H_{12}^{(1)}(5), C_{2n}$	-1.8	-1.9	-5.8	-6.0	-5.7	-9.8	-6.8

"The last two entries are computed by using the approximation of additivity.

bridged by the remaining hydrogen atoms. These probable or possible transient intermediates have been assigned these symmetries in theoretical studies.<sup>1</sup> However, a tendency for lowcoordinated boron to acquire additional interactions has previously been observed in theoretical studies in which electron correlation has been included<sup>2</sup> in a few other boranes. For that reason we have reexamined the stable geometries of  $B_3H_9$  and  $B_4H_{12}$  and show below that their most stable symmetries are  $C_{3v}$  and  $C_{2v}$ , respectively.

Optimized geometries were obtained by using stated symmetry constraints at the 3-21G level from the GAUSSIAN 82 program.<sup>3</sup> The same program was used for calculation of vibrational frequencies by the CPHF method<sup>3</sup> in a 3-21G basis. "Single-point" assumed geometries were used for subsequent calculations at the (Hartree-Fock) 6-31G level augmented by correlation at the MP2 and MP3 (Møller-Plesset) levels and polarization at the 6-31G\* level (d orbitals on B). Relative energies  $\Delta E = E_b - E_a$  for the distorted structure b as compared with those for the most symmetrical structure a were estimated by using the approximation<sup>4</sup>  $\Delta E(MP2/6-31G^*) \approx \Delta E(MP2/6-31G) +$ 

$$\Delta E(\text{HF}/6-31\text{G}^*) - \Delta E(\text{HF}/6-31\text{G})$$

The approximation that the correlation and polarization corrections are additive has been found to reproduce relative energies of these basis sets to within a few kcal/mol for isomerizations and several kcal/mol for more complex reactions.<sup>4</sup>

For  $B_3H_9$  the  $D_{3h}$  and  $C_{3v}$  structures, 1 and 2, respectively, in Figure 1, were optimized. The total energies (Table I) and relative

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- (4) We have approximated MP2/6-31G\* calculations by combining 6-31G\* and MP2 calculations and have approximated the MP3/6-31G\* calculations similarly. Additivity has been demonstrated to be accurate to within 1-4 kcal/mol of the full calculation. We suggest that in calculations which use additivity the appropriate basis set should be enclosed in brackets: McKee, M. L.; Lipscomb, W. N. J. Am, Chem. Soc. 1981, 103, 4673. Nobes, R. H.; Bouma, W. J.; Radom, L. Chem. Phys. Lett. 1982, 89, 497. McKee, M. L.; Lipscomb, W. N. Inorg. Chem. 1985, 24, 762.



**Figure 1.** Structures for  $B_3H_9$  and  $B_4H_{12}$ . In 5, + and - indicate displacements of 0.26 Å closer to or further from the plane containing the four bridges, as the symmetry lowers to  $C_{2\nu}$ .

**Table III.** Geometric Differences for  $B_3H_9$  and  $B_4H_{12}$  Optimized by Using the 3-21G Basis in Different Point Groups

	B−B, Å	В-Н <sub>ь</sub> , Å	B-H <sub>b</sub> -B, deg	B-B(trans), Å	vert dist, <sup>a</sup> Å
$1(D_{3h})$	2.351	1.305	128.5		0.0
<b>2</b> $(C_{3n})$	2.256	1.310	118.9		0.31
$3(D_{4h})$	2.606	1.306	172.8	3.686	0.0
$4(D_{2d})$	2.574	1.304	161.4	3.410	0.45 <sup>b</sup>
5 (C4.)	2.485	1.312	142.6	3.514	0.40
$5(C_{2v})$	2.471	1.310/1.310	141.3	3.235/3.593	0.12/0.64

<sup>a</sup>Distance between the plane containing the boron atoms and the plane containing the bridge hydrogens. <sup>b</sup>Vertical distance of boron atoms from the plane containing the bridge hydrogens.

energies (Table II) show that the  $C_{3v}$  structure is more stable at all levels of calculation. The difference is small, 2.2 kcal/mol at  $[MP3/6-31G^*]^4 + ZPC$  (zero-point vibrational corrections) and 4.4 kcal/mol if the ZPC corrections are omitted. This  $C_{3v}$ 

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Table IV. Calculated Vibrational Frequencies (cm<sup>-1</sup>) from 3-21G Optimized Geometries

$1 (D_{3h})$	<b>2</b> (C <sub>3v</sub> )	$3(D_{4h})$	$4(D_{2d})$	5 (C <sub>4v</sub> )	5 (C <sub>2v</sub> )
165i e"	218 e	44 <i>i</i> b <sub>1u</sub>	92 <i>i</i> b <sub>2</sub>	34 <i>i</i> b <sub>2</sub>	93 a <sub>1</sub>
214i a <sub>2</sub> "	274 a <sub>1</sub>	108 <i>i</i> b <sub>2u</sub>	143 <i>i</i> e	220 $a_1$	160 b <sub>1</sub>
430 a <sub>1</sub> '	491 a <sub>1</sub>	160i e''	146 <i>i</i> a <sub>2</sub>	242 e	208 a <sub>1</sub> /237 b <sub>2</sub>
491 e'	551 e	164 <i>i</i> a <sub>2u</sub>	72 a <sub>1</sub>	268 b <sub>1</sub>	239 a <sub>2</sub>
715 a <sub>2</sub> '	734 a <sub>2</sub>	278 b <sub>1g</sub>	212 b <sub>1</sub>	354 b <sub>1</sub>	309 a <sub>2</sub>
870 e''	862 e	361 b <sub>28</sub>	459 b <sub>2</sub>	396 b <sub>2</sub>	391 a <sub>1</sub>
970 a <sub>1</sub> ″	966 e	501 e'	508 a <sub>1</sub>	576 e	586 b <sub>2</sub> /592 b <sub>1</sub>
990 e'	1009 a <sub>2</sub>	504 a <sub>1g</sub>	531 e	577 a <sub>1</sub>	604 a <sub>1</sub>
1059 a <sub>2</sub> "	1165 a <sub>1</sub>	521 b <sub>1g</sub>	566 b <sub>1</sub>	617 a <sub>2</sub>	628 a <sub>2</sub>
1189 e'	1200 e	563 a <sub>2g</sub>	594 a <sub>2</sub>	634 b <sub>1</sub>	649 a <sub>2</sub>
1197 e''	1216 e	852 b <sub>1u</sub>	754 e	868 b <sub>2</sub>	808 b <sub>2</sub>
1224 a <sub>1</sub> '	1242 a <sub>1</sub>	879 e <sup>77</sup>	818 a <sub>1</sub>	886 e	$816 b_1/839 a_1$
1487 e <sup>7</sup>	1656 e	898 e'	972 b <sub>1</sub>	930 e	976 b <sub>1</sub> /985 b <sub>2</sub>
1592 a <sub>1</sub> '	1674 a <sub>1</sub>	902 a <sub>1u</sub>	992 e	987 a <sub>2</sub>	996 a <sub>2</sub>
2293 a <sub>2</sub> '	2228 a <sub>2</sub>	1120 e'	1132 e	1156 e	$1143 b_1/1153 b_2$
2402 e <sup>7</sup>	2278 e	1147 a <sub>2u</sub>	1146 b <sub>2</sub>	1171 a <sub>1</sub>	1178 a <sub>1</sub>
2752 e'	2764 e	1167 a <sub>1g</sub>	1180 a <sub>2</sub>	1212 b <sub>2</sub>	1208 a <sub>1</sub>
2762 a <sub>1</sub> '	2776 a <sub>1</sub>	1179 e <sup>//</sup>	1186 e	1243 b <sub>1</sub>	1230 a <sub>1</sub>
2872 e''	2885 e	1183 b <sub>2g</sub>	1194 b <sub>2</sub>	1258 e	$1241 b_2/1252 b_1$
2880 a <sub>2</sub> "	2896 a <sub>1</sub>	1189 b <sub>2u</sub>	1208 a <sub>1</sub>	1265 a <sub>1</sub>	1259 a <sub>1</sub>
		1243 b <sub>1g</sub>	1258 a <sub>1</sub>	1382 a <sub>1</sub>	1389 a <sub>1</sub>
		1246 e'	1264 e	1391 e	$1400 b_2/1404 a_2$
		1248 a <sub>1s</sub>	1272 b <sub>1</sub>	1 <b>39</b> 3 b <sub>1</sub>	1410 b <sub>1</sub>
		2183 a <sub>28</sub>	2240 a <sub>2</sub>	2110 a <sub>2</sub>	2143 a <sub>2</sub>
		2599 e <sup>7</sup> °	2574 e	2457 e	$2442 b_2/2474 b_1$
		2655 b <sub>2g</sub>	2636 b <sub>2</sub>	2509 b <sub>2</sub>	2508 a <sub>1</sub>
		2762 e <sup>7®</sup>	2761 e	2755 b <sub>2</sub>	2752 a <sub>1</sub>
		2763 a <sub>1e</sub>	2762 a <sub>1</sub>	2760 e	2753 b <sub>2</sub> /2762 b <sub>1</sub>
		2765 b <sub>2</sub>	2762 b <sub>2</sub>	2767 a <sub>1</sub>	2766 a <sub>1</sub>
		2868 b <sub>1</sub>	2871 a <sub>1</sub>	2873 b <sub>2</sub>	2871 a <sub>1</sub>
		2874 e''	2872 e	2883 e	2873 b <sub>2</sub> /2883 b <sub>1</sub>
		2882 a <sub>2u</sub>	2880 b <sub>2</sub>	2894 a <sub>1</sub>	2893 a <sub>1</sub>
		32	-	•	-

structure is more compact than the  $D_{3h}$  structure: the B-B distance is contracted from 2.351 to 2.256 Å, and the angle at the bridge hydrogen is reduced from 128.5 to 118.9° (Table III).

Vibrational frequencies (Table IV) for  $B_3H_9$  indicate that the  $C_{3v}$  structure is at a minimum, since all are positive and real. The three negative frequencies for the  $D_{3h}$  structure indicate that it is not a candidate for the transition state between interconverting  $C_{3v}$  structures. Although the negative frequency of 214*i* cm<sup>-1</sup> of the imaginary mode having  $a_2''$  symmetry could interconvert  $C_{3v}$  structures, the 65*i* cm<sup>-1</sup> frequency belongs to an imaginary mode in which two bridge hydrogens move in one direction and the third moves in the other direction out of the major plane of the molecule. Thus, the correct transition state is probably one in which one bridge "flips" first, followed by the other bridge hydrogens.

For  $B_4H_{12}$ , we find that the  $D_{2d}$ ,  $C_{4v}$ , and  $C_{2v}$  structures are all more stable than the  $D_{4h}$  structure at all levels of computation (Table III). The  $C_{2v}$  structure is more stable than the  $D_{4h}$  structure by 6.8 kcal/mol at the [MP3/6-31G\*] + ZPC level and is only 2.2 kcal/mol (3.5 kcal/mol) more stable than the  $C_{4v}$  ( $D_{2d}$ ) structure at this level (Table II). If ZPC corrections are omitted, the  $C_{2v}$  structure is 9.8 kcal/mol more stable than the  $D_{4h}$  structure and is 3.5 kcal/mol more stable than the  $D_{2d}$  structure at the [MP3/6-31G\*] level. Both of these lower symmetry structures are slightly more compact than the  $D_{4h}$  structure, as indicated by the shortening of B-B distances and decrease in the bond angle about the bridge hydrogens (Table III).

The vibrational frequencies (Table IV) reveal five imaginary modes for the  $D_{4h}$  structure, four imaginary modes for the  $D_{2d}$ structure, one imaginary mode for the  $C_{4v}$  structure, and no imaginary modes for the  $C_{2v}$  structure. Distortion of the  $C_{4v}$  structure to  $C_{2v}$  is indicated in 5, where diagonal borons are displaced 0.26 Å closer to (+) or further (-) form the plane containing the four bridge hydrogens while the average displacement is nearly identical with that found in the  $C_{4v}$  structure. Because of its negative eigenvalues, the  $D_{4h}$  structure cannot be the transition state for interconversion of  $C_{2v}$  structures, and we suggest that some sequence of consecutive or nonsimultaneous flips occurs during the interconversion such that the barrier remains below 6.8 kcal/mol. Both  $B_3H_9$  and  $B_4H_{12}$  are therefore double-minimum vibrational problems and should show the expected splitting of rotational lines.

Reduced symmetry<sup>2</sup> appears in BH<sub>5</sub>,  $B_2H_7^-$ ,  $B_4H_8$ , and  $B_5H_{11}$ and, here, in  $B_3H_9$  and  $B_4H_{12}$ . There are two main effects. The first is a tendency for low-coordinated boron to raise its coordination number, particularly if underuse of boron orbitals by bridge hydrogens occurs. The second is the conversion of terminal hydrogens to bridge hydrogens as long as strain is not excessive. This latter effect is primarily due to correlation: the pair of electrons in a bridge bond are separated by a nucleus, while these electrons show comparatively greater repulsion effects if they are in a single terminal bond.

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Registry No. B<sub>3</sub>H<sub>9</sub>, 36350-66-0; B<sub>4</sub>H<sub>12</sub>, 60349-62-4.

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# Vibrational Study of Ionic Association in Aprotic Solvents. 8. Copper(I) Thiocyanate in Soft Donor Solvents

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Copper(I) thiocyanate is insoluble in water and in most usual solvents, but it is readily soluble in N,N-dimethylthioformamide (DMTF), which is highly polar, donating through the soft S atom. This salt is also slightly soluble  $((\sim 1-4) \times 10^{-2} \text{ M})$  in tetra-hydrothiophene (THT), dimethyl sulfoxide (Me<sub>2</sub>SO), and hexamethylphosphoric triamide (HMPT). The existing species and the bonding mode of SCN can be identified by infrared spectroscopy.

## **Experimental Section**

THT was distilled on a 20-plate column, and DMTF was distilled under reduced pressure. Other solvents were of the best commercial grade. All solvents were stored on 4-Å molecular sieves. Solutions were prepared in a drybox. IR spectra were recorded on an interferometric Bruker IFS 45 spectrometer and Raman spectra on a Coderg T 800 spectrometer.

# **Results and Discussion**

The CuSCN solutions in pure DMTF were first investigated in the  $\nu(CN)$  and  $\nu(CS)$  regions of SCN. In both regions two absorption bands were observed at 2055 and 2076 cm<sup>-1</sup> (Figure 1b) and at 737 and 780 cm<sup>-1</sup>. The bands at 2055 and 737 cm<sup>-1</sup> are characteristic of the SCN<sup>-</sup> ion. They appeared at the same frequency in spectra of NBu<sub>4</sub>SCN solutions (Figure 1a). This assignment was confirmed by dilution effects on the relative intensities in the v(CN) region. The other two bands were assigned to the isothiocyanate ion pair CuNCS because the frequency shift of  $\nu(CS)$  relative to SCN<sup>-</sup> is positive while it is negative for thiocyanate complexes such as AgSCN.<sup>1,2</sup> The frequency shifts (27 and 43 cm<sup>-1</sup>) of  $\nu$ (CN) and  $\nu$ (CS) are in the same ratio as for LiNCS (14 and 30 cm<sup>-1</sup>) and somewhat larger, which indicates that the structures of both ion pairs are similar. The isothiocyanate structure of CuNCS was also proved by the variation in integrated intensity between SCN<sup>-</sup> and CuNCS. This variation was more easily measured in THT where the integrated intensity of  $\nu(CN)$ rises from 17100 in NBu<sub>4</sub>SCN to 20700 mol<sup>-1</sup> L cm<sup>-2</sup> in CuNCS.

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